

Heat equation in n dimensions:

$$c\rho \frac{\partial u}{\partial t} - \nabla \cdot (k_0 \nabla u) = Q$$

Initial conditions: obvious $u(x, 0) = u_0(x)$
for $x \in \mathbb{R}$

Boundary conditions:

Heat bath: $u(x, t) = g(x)$ for $x \in \partial \mathbb{R}$

Insulated: $\varphi(x, t) \cdot \hat{n}(x) = 0$ for $x \in \partial \mathbb{R}$

Fourier's law:

$$\varphi(x, t) = -k_0(x) \nabla u(x, t)$$

$$\nabla u(x, t) \cdot \hat{n}(x) = 0 \quad \text{for } x \in \partial \mathbb{R}.$$

Newton's law of cooling:

$$\varphi(x, t) \cdot \hat{n}(x) = H(u(x, t) - g(x, t)) \quad \text{for } x \in \partial \mathbb{R}.$$

Predict the future: Equilibrium solution as $t \rightarrow \infty \dots$

then u doesn't change in time as it's in equilibrium. So

$$\frac{\partial u}{\partial t} = 0$$

Heat equation

$$c\rho \frac{\partial u}{\partial t} - \nabla \cdot (k_0 \nabla u) = Q$$

before crossing out $\frac{\partial u}{\partial t}$ this was a parabolic problem.

$$-\nabla \cdot (k_0 \nabla u) = Q$$

Simplifying assumption: The conductivity is constant throughout the region R . $k_0(x) = k_0$

$$-k_0 \nabla \cdot \nabla u = Q$$

Poisson equation

$$\nabla^2 u = -\frac{Q(x)}{k_0}$$

time independent Q

unlike the 1-D case this is not an ODE but still a PDE

This is an elliptic problem ... i.e. statics problem equilibrium state

Set $Q=0$ then

$$\nabla^2 u = 0$$

Laplace equation

Laplace operator ... in some books $\nabla^2 = \Delta$

Chapter 2 Separation of Variables ...

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2}$$

simplest 1D heat equation

I.C. $u(x, 0) = u_0(x)$ for $x \in [0, L]$

B.C. $u(x,t) \approx 0$ for $x=0$ and $x=L$

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2}$$

This is a linear PDE...

In linear algebra we had matrices. What was a matrix? A linear function...

In PDE we have derivatives. What is a derivative? A linear operation on functions...

works like the matrix M

$$\rightarrow h(u) = \frac{\partial u}{\partial t} - k \cdot \frac{\partial^2 u}{\partial x^2}$$

$$h(c_1 u_1 + c_2 u_2) = \frac{\partial}{\partial t} (c_1 u_1 + c_2 u_2) - k \frac{\partial^2}{\partial x^2} (c_1 u_1 + c_2 u_2)$$

heat distributions

const

const

$$= c_1 \frac{\partial u_1}{\partial t} + c_2 \frac{\partial u_2}{\partial t} - c_1 k \frac{\partial^2 u_1}{\partial x^2} - c_2 k \frac{\partial^2 u_2}{\partial x^2}$$

$$= c_1 h(u_1) + c_2 h(u_2)$$

Thus $h(c_1 u_1 + c_2 u_2) = c_1 h(u_1) + c_2 h(u_2)$

That means h is a linear operator on the space of heat distributions.

In linear algebra we solve

$$Ax = b$$

In PDEs we solve

$$L(u) = \frac{Q}{c\rho}$$

← both are linear problems.

Linear algebra

Consider the Homogeneous problem $Ax = 0$

Solutions were called $\text{Nul}(A) = \{x : Ax = 0\}$

PDEs

Consider the Homogeneous problem $Lu = 0$

Solutions are the $\text{Ker}(L) = \{u : Lu = 0\}$.

Separation of variables

Assume $u(x, t) = f(x)G(t)$

and plug it in...

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2}$$

$$\frac{\partial u}{\partial t} = \frac{\partial}{\partial t} (f(x)G(t)) = f(x) \frac{d}{dt} G(t)$$

$$\frac{\partial^2 u}{\partial x^2} = \frac{\partial^2}{\partial x^2} (f(x)G(t)) = G(t) \frac{d^2}{dx^2} f(x)$$

Magic

$$\varphi(x) \frac{d}{dt} G(t) = k G(t) \frac{d^2}{dx^2} \varphi(x)$$

$$\frac{1}{k G(t)} \frac{d}{dt} G(t) = \frac{1}{\varphi(x)} \frac{d^2}{dx^2} \varphi(x) = -\lambda$$

since the left doesn't depend on x
and the right doesn't depend on t
then since they are equal ...
they don't depend on either x or t

Therefore

$$\frac{1}{k G(t)} \frac{dG(t)}{dt} = -\lambda$$

$$\frac{dG}{dt} = -\lambda k G$$

and

$$\frac{1}{\varphi(x)} \frac{d^2}{dx^2} \varphi(x) = -\lambda$$

$$\frac{d^2}{dx^2} \varphi = -\lambda \varphi$$

these are like the
eigenvalue-eigenvector
equation $Ax = \lambda x$
in linear algebra...

These eigenvalue problems involving differential operators
are ODE's and I already know how to solve them

Solve this ODE ...

$$\frac{dG}{dt} = -\lambda k G$$

Guess method $G_2(t) = G(0) e^{-\lambda k t}$

Separation of Vbls...

$$\frac{dG}{dt} = -\lambda k G$$

$$\int \frac{dG}{G} = \int -\lambda k dt$$

$$\ln G = -\lambda k t + C$$

$$G = e^{\ln G} = e^{-\lambda k t + C} = e^{-\lambda k t} e^C$$

Thus

$$G(t) = G(0) e^{-\lambda k t}$$

Solve the other ODE:

$$\frac{d^2}{dx^2} \phi = -\lambda \phi$$

B.C. $u(x,t) = 0$ for $x=0$ and $x=L$

$$u(x,t) = \phi(x) G(t) = 0$$

for $x=0, x=L$
and all times

Boundary value
ODE

Then $f(0) = 0$ and $f(L) = 0$

$$\frac{d^2 f}{dx^2} = -\lambda f \quad \text{such that } f(0) = 0 \text{ and } f(L) = 0$$

$$f(x) = A \cos(\sqrt{\lambda} x) + B \sin(\sqrt{\lambda} x)$$

$$f(0) = 0 = A = 0$$

$$f(L) = 0 = B \sin(\sqrt{\lambda} L) = 0$$

set of
integers
↓

thus $\sqrt{\lambda} L = n\pi$ for some $n \in \mathbb{Z}$

Solve for the eigenvalues...

$$\sqrt{\lambda} = \frac{n\pi}{L}$$

$$\lambda = \left(\frac{n\pi}{L}\right)^2$$

To solve for B we will use initial condition

$$u(x, 0) = u_0(x)$$

Next time ...