

$$\frac{d}{dx} \left( p(x) \frac{d\phi}{dx} \right) + q(x)\phi + \lambda\sigma(x)\phi = 0,$$

$$a < x < b,$$

$$\beta_1\phi(a) + \beta_2 \frac{d\phi}{dx}(a) = 0$$

$$\beta_3\phi(b) + \beta_4 \frac{d\phi}{dx}(b) = 0,$$

$p > 0$  and  $\sigma > 0$  everywhere

1. All the eigenvalues  $\lambda$  are real.

2. There exist an infinite number of eigenvalues:

$$\lambda_1 < \lambda_2 < \dots < \lambda_n < \lambda_{n+1} < \dots$$

a. There is a smallest eigenvalue, usually denoted  $\lambda_1$ .

b. There is not a largest eigenvalue and  $\lambda_n \rightarrow \infty$  as  $n \rightarrow \infty$ .

3. Corresponding to each eigenvalue  $\lambda_n$ , there is an eigenfunction, denoted  $\phi_n(x)$  (which is unique to within an arbitrary multiplicative constant).  $\phi_n(x)$  has exactly  $n - 1$  zeros for  $a < x < b$ .

4. The eigenfunctions  $\phi_n(x)$  form a "complete" set, meaning that any piecewise smooth function  $f(x)$  can be represented by a generalized Fourier series of the eigenfunctions:

$$f(x) \sim \sum_{n=1}^{\infty} a_n \phi_n(x).$$

Furthermore, this infinite series converges to  $[f(x+) + f(x-)]/2$  for  $a < x < b$  (if the coefficients  $a_n$  are properly chosen).

5. Eigenfunctions belonging to different eigenvalues are orthogonal relative to the weight function  $\sigma(x)$ . In other words,

$$\int_a^b \phi_n(x)\phi_m(x)\sigma(x) dx = 0 \quad \text{if } \lambda_n \neq \lambda_m.$$

6. Any eigenvalue can be related to its eigenfunction by the **Rayleigh quotient**:

$$\lambda = \frac{-p\phi \, d\phi/dx|_a^b + \int_a^b [p(d\phi/dx)^2 - q\phi^2] dx}{\int_a^b \phi^2 \sigma dx},$$

where the boundary conditions may somewhat simplify this expression.

PDE: 
$$c\rho \frac{\partial u}{\partial t} = \frac{\partial}{\partial x} \left( K_0 \frac{\partial u}{\partial x} \right)$$

$c = c(x)$   
 $\rho = \rho(x)$   
 $K_0 = K_0(x)$

BC: 
$$\begin{aligned} u(0, t) &= 0 \\ \frac{\partial u}{\partial x}(L, t) &= 0 \end{aligned}$$

IC: 
$$u(x, 0) = f(x).$$

Idea: separation of variables  
 superposition  
 orthogonality

Let  $u(x, t) = \phi(x)h(t)$   
 and plug it in...

$$c(x)\rho(x)\phi(x)h'(t) = \frac{\partial}{\partial x} \left( K_0(x)\phi'(x)h(t) \right)$$

$$c(x)\rho(x)\phi(x)h'(t) = h(t) \frac{d}{dx} (K_0(x)\phi'(x))$$

$$\frac{h'(t)}{h(t)} = \frac{\frac{d}{dx} (K_0(x)\phi'(x))}{c(x)\rho(x)\phi(x)} = -\lambda$$

Get two ODEs...

$$h'(t) = -\lambda h(t)$$

✓ solve

$$h(t) = c e^{-\lambda t}$$

$\uparrow$   
 $a_n$

$$\frac{d}{dx} (K_0(x)\phi'(x)) + \lambda c(x)\rho(x)\phi(x) = 0$$

$$\frac{d}{dx} \left( p(x) \frac{d\phi}{dx} \right) + q(x)\phi + \lambda \sigma(x)\phi = 0,$$

note also  $p > 0$  and  $q \geq 0$

apply theorem about Sturm-Liouville problems here...

$$\begin{aligned} u(0, t) &= 0 \\ \frac{\partial u}{\partial x}(L, t) &= 0 \end{aligned}$$

$$\begin{aligned} \beta_1 \phi(a) + \beta_2 \frac{d\phi}{dx}(a) &= 0 \\ \beta_3 \phi(b) + \beta_4 \frac{d\phi}{dx}(b) &= 0, \end{aligned}$$

Check boundary conditions  $a=0$ ,  $b=L$

$$\beta_1=1, \beta_2=0, \beta_3=0, \beta_4$$

shows they are of the form needed by the Theorem,

by (2) and (3) there are an infinite # of eigenvalues  $\lambda_n$  with corresponding eigen functions  $\phi_n$  such that

$$\frac{d}{dx} (k_0(x) \phi_n'(x)) + \lambda_n c(x) \rho(x) \phi_n(x) = 0$$

That is this ODE has lots of solutions...

Superposition principle

$$u(x, t) = \sum_{n=1}^{\infty} a_n \phi_n(x) e^{-\lambda_n t}$$

Note that  $u(x, t)$  satisfies the boundary conditions because they were homogeneous and because each of the  $\phi_n$ 's satisfy them.

Now I want to solve for the initial condition

$$u(x, 0) = f(x).$$

Thus

$$u(x,0) = \sum_{n=1}^{\infty} a_n \phi_n(x) = f(x)$$

Use orthogonality <sup>⑤</sup> to solve for the  $a_n$ 's under the assumption that  $f(x)$  can even be expressed this way.

④ piecewise smooth function  $f(x)$  can be represented by a generalized Fourier series of the eigenfunctions:

$$f(x) \sim \sum_{n=1}^{\infty} a_n \phi_n(x).$$

Furthermore, this infinite series converges to  $[f(x+) + f(x-)]/2$

So ④ in the theorem says  $f(x)$  can be written this way provided it's piecewise smooth.

$$\int_a^b \phi_n(x) \phi_m(x) \sigma(x) dx = 0 \quad \text{if } \lambda_n \neq \lambda_m.$$

Thus

$$\int_0^L \phi_n(x) \phi_m(x) c(x) p(x) dx = 0 \quad \text{if } \lambda_n \neq \lambda_m$$

Use orthogonality to solve for the  $a_n$ 's so that

$$\sum_{n=1}^{\infty} a_n \phi_n(x) = f(x)$$

$$\sum_{n=1}^{\infty} a_n \phi_n(x) \phi_m(x) c(x) p(x) = f(x) \phi_m(x) c(x) p(x)$$

$$\int_0^L \sum_{n=1}^{\infty} a_n \phi_n(x) \phi_m(x) c(x) p(x) dx = \int_0^L f(x) \phi_m(x) c(x) p(x) dx$$

$$\sum_{n=1}^{\infty} a_n \int_0^L \phi_n(x) \phi_m(x) c(x) \rho(x) dx = \int_0^L f(x) \phi_m(x) c(x) \rho(x) dx$$

non zero only when  $n=m$ .

If  $n=m$

$$\int_0^L \phi_n(x) \phi_m(x) c(x) \rho(x) dx = \int_0^L \phi_m(x)^2 c(x) \rho(x) dx$$

$$a_m \int_0^L \phi_m(x)^2 c(x) \rho(x) dx = \int_0^L f(x) \phi_m(x) c(x) \rho(x) dx$$

Solution is

$$u(x,t) = \sum_{n=1}^{\infty} a_n \phi_n(x) e^{-\lambda_n t}$$

where

$$a_m = \frac{\int_0^L f(x) \phi_m(x) c(x) \rho(x) dx}{\int_0^L \phi_m(x)^2 c(x) \rho(x) dx}$$

What about ① and ⑥ in the theorem? Useful for interpretation...  $\lambda_n$  are real.

Any eigenvalue can be related to its eigenfunction by the Rayleigh quotient:

$$\lambda_n = \frac{-p \phi_n' \phi_n / dx \Big|_0^L + \int_0^L [p (d\phi_n/dx)^2 - q \phi_n^2] dx}{\int_0^L \phi_n^2 \sigma dx}$$

where the boundary conditions may somewhat simplify this expression.

Boundary conditions

$$\begin{cases} u(0, t) = 0 \\ \frac{\partial u}{\partial x}(L, t) = 0 \end{cases}$$

$\phi_n(0) = 0$

$\phi_n'(L) = 0$

$$-k_0 \phi_n \phi_n' \Big|_0^L = -k_0(L) \phi_n(L) \phi_n'(L) + k_0(0) \phi_n(0) \phi_n'(0) = 0$$

$$\lambda_n = \frac{\int_0^L k_0(x) \phi_n'(x)^2 dx}{\int_0^L \phi_n(x)^2 c(x) \rho(x) dx}$$

Since  $\phi_n(x) = 0$  is never an eigenfunction then

$$\int_0^L \phi_n(x)^2 c(x) \rho(x) dx > 0$$

What if  $\phi_n'(x) = 0$ ? Could it happen?

If it did that means  $\phi_n(x) = \text{const.}$

$\phi_n(0) = 0 \leftarrow \text{Boundary cond.} \dots$

Then  $\phi_n(x) = 0 \dots$  which is not eigenfunction  $\dots$

Conclusion  $\phi_n'(x) \neq 0$ .

Thus

$$\int_0^L k_0(x) \phi_n'(x)^2 dx > 0$$

positive square of non-zero

Therefore

$$\lambda_n = \frac{\int_0^L k_0(x) \phi_n'(x)^2 dx}{\int_0^L \phi_n(x)^2 c(x) \rho(x) dx} > 0$$

Thus every term in the solution

$$u(x,t) = \sum_{n=1}^{\infty} a_n \phi_n(x) e^{-\lambda_n t}$$

↑ exponentially decays to zero

Qualitatively the behavior as  $t \rightarrow \infty$  is similar to the constant coefficient heat equation with same boundary conditions...

More stuff about Rayleigh quotient...

Linear algebra case... Spectral theorem for symmetric Matrices...

If  $A = A^T$  then there is an orthonormal basis of eigenvectors  $x_n$  with real eigen values  $\lambda_n$  such that

$$Ax_n = \lambda_n x_n \quad \text{and} \quad x_n \cdot x_m = \begin{cases} 1 & \text{if } n=m \\ 0 & \text{if } n \neq m \end{cases}$$

Recall Rayleigh Quotient

$$R(x) = \frac{Ax \cdot x}{x \cdot x}$$

Note: if  $x = x_n$  then  $R(x_n) = \lambda_n$

$$Ax \cdot x = \lambda x \cdot x$$

If  $x \neq x_n$  then since they form a basis

$$x = c_1 x_1 + c_2 x_2 + \dots + c_n x_n$$

$$R(x) = \frac{A(c_1 x_1 + c_2 x_2 + \dots + c_n x_n) \cdot (c_1 x_1 + c_2 x_2 + \dots + c_n x_n)}{(c_1 x_1 + c_2 x_2 + \dots + c_n x_n) \cdot (c_1 x_1 + c_2 x_2 + \dots + c_n x_n)}$$

$$R(x) = \frac{(c_1 \lambda_1 x_1 + c_2 \lambda_2 x_2 + \dots + c_n \lambda_n x_n) \cdot (c_1 x_1 + c_2 x_2 + \dots + c_n x_n)}{(c_1 x_1 + c_2 x_2 + \dots + c_n x_n) \cdot (c_1 x_1 + c_2 x_2 + \dots + c_n x_n)}$$

$$= \frac{\lambda_1 c_1^2 + \lambda_2 c_2^2 + \dots + \lambda_n c_n^2}{c_1^2 + c_2^2 + \dots + c_n^2} \quad \leftarrow \text{this is a weighted average of the } \lambda_i\text{'s.}$$

$$\min \{ \lambda_i \} \leq \frac{\lambda_1 c_1^2 + \lambda_2 c_2^2 + \dots + \lambda_n c_n^2}{c_1^2 + c_2^2 + \dots + c_n^2} \leq \max \{ \lambda_i \}$$

Order the eigenvalues

$$\lambda_1 \leq \lambda_2 \leq \lambda_3 \leq \dots \leq \lambda_n$$

then

$$\lambda_1 = \min \{ R(x) : x \in \mathbb{R}^n \}$$

$$\lambda_n = \max \{ R(x) : x \in \mathbb{R}^n \}.$$

For Sturm-Liouville problems the same works, except not for the max eigenvalue because they go to  $\infty$ .

$$\lambda_1 \leq RQ[u_T] = \frac{-p u_T du_T/dx|_a^b + \int_a^b [p (du_T/dx)^2 - q u_T^2] dx}{\int_a^b u_T^2 \sigma dx} \quad (5.6.6)$$

smallest eigenvalue..

bounded above by the Rayleigh quotient of any function that satisfies the boundary condition.



One gets exact equality by minimizing...

$$\lambda_1 = \min \frac{-pu \, du/dx|_a^b + \int_a^b [p (du/dx)^2 - qu^2] \, dx}{\int_a^b u^2 \sigma \, dx},$$

Back to PDE: knowing  $\lambda_1$  tells how fast

$$u(x,t) = \sum_{n=1}^{\infty} a_n \phi_n(x) e^{-\lambda_n t}$$

↑ exponentially  
decays to zero

decays because  $\lambda_1$  is the slowest decaying term in the series.