

20- Apr - 2023

~~[16 Feb 2023]~~ Quiz 2

April 20

Quiz 2 will be given in class on ~~February 16~~. Please be prepared to solve the Laplace equation as in Exercise 2.5.1 from the text and the wave equation as in Exercises 4.4.2 and 4.4.3.

**2.5.1.** Solve Laplace's equation inside a rectangle  $0 \leq x \leq L$ ,  $0 \leq y \leq H$ , with the following boundary conditions [*Hint*: Separate variables. If there are two homogeneous boundary conditions in  $y$ , let  $u(x, y) = h(x)\phi(y)$ , and if there are two homogeneous boundary conditions in  $x$ , let  $u(x, y) = \phi(x)h(y)$ .]:

**\*(a)**  $\frac{\partial u}{\partial x}(0, y) = 0$ ,  $\frac{\partial u}{\partial x}(L, y) = 0$ ,  $u(x, 0) = 0$ ,  $u(x, H) = f(x)$

**(b)**  $\frac{\partial u}{\partial x}(0, y) = g(y)$ ,  $\frac{\partial u}{\partial x}(L, y) = 0$ ,  $u(x, 0) = 0$ ,  $u(x, H) = 0$

**\*(c)**  $\frac{\partial u}{\partial x}(0, y) = 0$ ,  $u(L, y) = g(y)$ ,  $u(x, 0) = 0$ ,  $u(x, H) = 0$

**(d)**  $u(0, y) = g(y)$ ,  $u(L, y) = 0$ ,  $\frac{\partial u}{\partial y}(x, 0) = 0$ ,  $u(x, H) = 0$

**\*(e)**  $u(0, y) = 0$ ,  $u(L, y) = 0$ ,  $u(x, 0) - \frac{\partial u}{\partial y}(x, 0) = 0$ ,  $u(x, H) = f(x)$

**(f)**  $u(0, y) = f(y)$ ,  $u(L, y) = 0$ ,  $\frac{\partial u}{\partial y}(x, 0) = 0$ ,  $\frac{\partial u}{\partial y}(x, H) = 0$

**(g)**  $\frac{\partial u}{\partial x}(0, y) = 0$ ,  $\frac{\partial u}{\partial x}(L, y) = 0$ ,  $u(x, 0) = \begin{cases} 0 & x > L/2 \\ 1 & x < L/2 \end{cases}$ ,  $\frac{\partial u}{\partial y}(x, H) = 0$

**(h)**  $u(0, y) = 0$ ,  $u(L, y) = g(y)$ ,  $u(x, 0) = 0$ ,  $u(x, H) = 0$

**4.4.2.** In Section 4.2 it was shown that the displacement  $u$  of a nonuniform string satisfies

$$\rho_0 \frac{\partial^2 u}{\partial t^2} = T_0 \frac{\partial^2 u}{\partial x^2} + Q,$$

where  $Q$  represents the vertical component of the body force per unit length. If  $Q = 0$ , the partial differential equation is homogeneous. A slightly different homogeneous equation occurs if  $Q = \alpha u$ .

- (a) Show that if  $\alpha < 0$ , the body force is restoring (toward  $u = 0$ ). Show that if  $\alpha > 0$ , the body force tends to push the string further away from its unperturbed position  $u = 0$ .
- (b) Separate variables if  $\rho_0(x)$  and  $\alpha(x)$  but  $T_0$  is constant for physical reasons. Analyze the time-dependent ordinary differential equation.
- \* (c) Specialize part (b) to the constant coefficient case. Solve the initial value problem if  $\alpha < 0$ :

$$\begin{aligned} u(0, t) &= 0, & u(x, 0) &= 0 \\ u(L, t) &= 0, & \frac{\partial u}{\partial t}(x, 0) &= f(x). \end{aligned}$$

What are the frequencies of vibration?

**4.4.3.** Consider a slightly damped vibrating string that satisfies

$$\rho_0 \frac{\partial^2 u}{\partial t^2} = T_0 \frac{\partial^2 u}{\partial x^2} - \beta \frac{\partial u}{\partial t}.$$

- (a) Briefly explain why  $\beta > 0$ .
- \* (b) Determine the solution (by separation of variables) that satisfies the boundary conditions

$$u(0, t) = 0 \quad \text{and} \quad u(L, t) = 0$$

and the initial conditions

$$u(x, 0) = f(x) \quad \text{and} \quad \frac{\partial u}{\partial t}(x, 0) = g(x).$$

You can assume that this frictional coefficient  $\beta$  is relatively small ( $\beta^2 < 4\pi^2 \rho_0 T_0 / L^2$ ).

\* (e)  $u(0, y) = 0$ ,  $u(L, y) = 0$ ,  $u(x, 0) - \frac{\partial u}{\partial y}(x, 0) = 0$ ,  $u(x, H) = \underline{\underline{f(x)}}$  not homogeneous

$$u(x, y) = \varphi(x) h(y)$$

Laplace equation  $u_{xx} + u_{yy} = 0$  or  $\Delta u = 0$

$$\varphi''(x) h(y) + \varphi(x) h''(y) = 0$$

$$\frac{\varphi''(x)}{\varphi(x)} = - \frac{h''(y)}{h(y)} = -\lambda$$

Thus,

$$\varphi''(x) = -\lambda \varphi(x)$$

$$\varphi(0) = 0, \quad \varphi(L) = 0$$

and  $h''(y) = \lambda h(y)$

$$h(0) - h'(0) = 0, \quad h(H) = C$$

Solve for  $\varphi$ :

$$\varphi(x) = a \cos \sqrt{\lambda} x + b \sin \sqrt{\lambda} x$$

$$\varphi(0) = a = 0 \quad \text{so } a = 0$$

$$\varphi(L) = b \sin \sqrt{\lambda} L = 0$$

Since  $b \neq 0$  then  $\sqrt{\lambda} L = n\pi$

$$\sqrt{\lambda} = \frac{n\pi}{L}$$

$$\lambda = \left(\frac{n\pi}{L}\right)^2$$

Next solve for  $h \dots$

$$h(y) = \alpha e^{r\sqrt{\lambda} y} + \beta e^{-r\sqrt{\lambda} y}$$

General idea  $\dots$  substitute  $e^{rt}$  to obtain

$$r^2 e^{rt} = \lambda e^{rt}$$

$$r^2 = \lambda \quad \text{so} \quad r = \pm \sqrt{\lambda}$$

relationship for  $\alpha$  and  $\beta$  so that

$$h(0) - h'(0) = 0$$

$$h'(y) = \alpha \sqrt{\lambda} e^{r\sqrt{\lambda} y} - \beta \sqrt{\lambda} e^{-r\sqrt{\lambda} y}$$

$$h(0) - h'(0) = \alpha + \beta - \alpha \sqrt{\lambda} + \beta \sqrt{\lambda} = 0$$

$$\alpha (1 + \sqrt{\lambda}) - \beta (1 - \sqrt{\lambda}) = 0$$

$$\alpha = \beta \left( \frac{1 - \sqrt{\lambda}}{\sqrt{\lambda} + 1} \right)$$

Therefore by superposition

$$u(x,t) = \sum_{n=1}^{\infty} \beta_n \sin \frac{n\pi x}{L} \left[ \left( \frac{1-\sqrt{\lambda}}{1+\sqrt{\lambda}} \right) e^{\frac{n\pi}{L}y} + e^{-\frac{n\pi}{L}y} \right]$$

Panic... should have use hyperbolic functions, maybe...  
 Quick check if better... yes worth checking... (i)

Next solve for h...

$$h(y) = \alpha e^{\sqrt{\lambda}y} + \beta e^{-\sqrt{\lambda}y}$$

$$h(y) = \gamma \cosh \sqrt{\lambda}y + \delta \sinh \sqrt{\lambda}y$$

$$h'(y) = \sqrt{\lambda} \gamma \sinh \sqrt{\lambda}y + \sqrt{\lambda} \delta \cosh \sqrt{\lambda}y$$

$$h(0) - h'(0) = \gamma - \sqrt{\lambda} \delta = 0 \quad \gamma = \sqrt{\lambda} \delta \quad \cosh \theta = \frac{e^{\theta} + e^{-\theta}}{2}$$

Therefore by superposition

$$\sinh \theta = \frac{e^{\theta} - e^{-\theta}}{2}$$

$$u(x,t) = \sum_{n=1}^{\infty} \delta_n \sin \frac{n\pi x}{L} \left[ \sqrt{\lambda} \cosh \frac{n\pi y}{L} + \sinh \frac{n\pi y}{L} \right]$$

$$= \sum_{n=1}^{\infty} \delta_n \sin \frac{n\pi x}{L} e^{\frac{n\pi y}{L}}$$

Find the coefficient  $\delta_n$  to satisfy the remaining boundary condition

$$u(x,H) = f(x)$$

$$u(x,H) = \sum_{n=1}^{\infty} \delta_n \sin \frac{n\pi x}{L} \left[ \underbrace{\left[ \frac{n\pi}{L} \cosh \frac{n\pi H}{L} + \sinh \frac{n\pi H}{L} \right]}_{\text{constant}} \right] = f(x)$$

$$\int_0^L \sum_{n=1}^{\infty} \delta_n \sin \frac{n\pi x}{L} \left[ \int_0^L \cosh \frac{n\pi H}{L} + \sinh \frac{n\pi H}{L} \right] \sin \frac{n\pi x}{L} dx = \int_0^L f(x) \sin \frac{n\pi x}{L} dx$$

$$\int_0^L \delta_n \left[ \int_0^L \cosh \frac{n\pi H}{L} + \sinh \frac{n\pi H}{L} \right] \sin \frac{n\pi x}{L} dx = \int_0^L f(x) \sin \frac{n\pi x}{L} dx$$

$$\delta_n = \frac{1}{\int_0^L \left[ \cosh \frac{n\pi H}{L} + \sinh \frac{n\pi H}{L} \right] \sin \frac{n\pi x}{L} dx} \int_0^L f(x) \sin \frac{n\pi x}{L} dx$$

Solution

$$u(x, t) = \sum_{n=1}^{\infty} \delta_n \sin \frac{n\pi x}{L} \left[ \cosh \frac{n\pi y}{L} + \sinh \frac{n\pi y}{L} \right]$$

where

$$\delta_n = \frac{1}{\int_0^L \left[ \cosh \frac{n\pi H}{L} + \sinh \frac{n\pi H}{L} \right] \sin \frac{n\pi x}{L} dx} \int_0^L f(x) \sin \frac{n\pi x}{L} dx$$

## CHAPTER 12

# The Method of Characteristics for Linear and Quasilinear Wave Equations

$$\frac{\partial^2 u}{\partial t^2} - c^2 \frac{\partial^2 u}{\partial x^2} = 0.$$

s that it can be "factored" in two ways:

$$\left( \frac{\partial}{\partial t} + c \frac{\partial}{\partial x} \right) \left( \frac{\partial u}{\partial t} - c \frac{\partial u}{\partial x} \right) = 0$$

$$\left( \frac{\partial}{\partial t} - c \frac{\partial}{\partial x} \right) \left( \frac{\partial u}{\partial t} + c \frac{\partial u}{\partial x} \right) = 0.$$

erivative terms vanish in both. If we let

$$w = \frac{\partial u}{\partial t} - c \frac{\partial u}{\partial x}$$

$$v = \frac{\partial u}{\partial t} + c \frac{\partial u}{\partial x}$$

$$\begin{cases} \left( \frac{\partial}{\partial t} + c \frac{\partial}{\partial x} \right) w = 0 \\ \left( \frac{\partial}{\partial t} - c \frac{\partial}{\partial x} \right) v = 0 \end{cases}$$

$$\left( \frac{\partial}{\partial t} + c \frac{\partial}{\partial x} \right) w = 0$$

$$\frac{\partial w}{\partial t} + c \frac{\partial w}{\partial x} = 0$$

$$\frac{\partial v}{\partial t} - c \frac{\partial v}{\partial x} = 0$$

How to solve  $\frac{\partial w}{\partial t} + c \frac{\partial w}{\partial x} = 0$  for the function  $w(x,t)$ ?

parameterize  $x$  as a function of  $t$

or  $t$  ~~as~~  $x$  to get an ODE

In this  $x = x(t)$  and consider  $w(x(t), t)$  which is now only a function of  $t$ .

$$\frac{d}{dt} w(x(t), t) = w_x(x(t), t) x'(t) + w_t(x(t), t)$$

$$\frac{\partial w}{\partial t} + c \frac{\partial w}{\partial x} = 0$$

If  $x'(t) = c$  then  $\frac{d}{dt} w(x(t), t) = 0$

First solve the ODE... to get the characteristic directions

Then solve this ode along those directions

$x(t) = ct + x(0)$  *plug in*

$$\frac{d}{dt} w(ct + x(0), t) = 0$$

Thus  $w(ct + x(0), t) = \text{const} = w(x(0), 0)$  *at t=0*

Solving this... but what was the initial condition

$$\frac{\partial w}{\partial t} + c \frac{\partial w}{\partial x} = 0$$

$$w(x, 0) = f(x)$$

Thus  $w(ct + x, t) = f(x)$

This is the initial cond in terms of the solution. Want solution in terms of initial condition

if  $\alpha = ct + x$   
 then  $x = \alpha - ct$

$$w(\alpha, t) = f(\alpha - ct)$$

Solution in terms of the initial condition...

ordinary d.e.

$$\frac{d}{dt} w(x(t), t) = w_x(x(t), t) x'(t) + w_t(x(t), t) = 0$$

$$\frac{\partial w}{\partial t} + c \frac{\partial w}{\partial x} = 0$$

Suppose  $t$  as a funct of  $x$

$w(x, t(x))$  is a function of  $x$

$$\frac{d}{dx} w(x, t(x)) = \frac{\partial w}{\partial x} \frac{dx}{dx} + \frac{\partial w}{\partial t} \frac{dt}{dx} = w_x + w_t t'(x)$$

$$t'(x) = \frac{1}{c}$$