

What's next?

Quiz 2 on April 19

HW 5 due today.

HW 6 is extra and the last one... it will be used to replace any lower HW score if you turn it in.

PDE: $\rho \frac{\partial^2 u}{\partial t^2} = T_0 \frac{\partial^2 u}{\partial x^2}$ BC: $\begin{matrix} u(0, t) = 0 \\ u(L, t) = 0 \end{matrix}$ IC: $\begin{matrix} u(x, 0) = f(x) \\ \frac{\partial u}{\partial t}(x, 0) = g(x). \end{matrix}$

$\rho = \rho(x)$ so the density is no longer constant.

Use separation of variables and superposition to solve the PDE.

$$u(x, t) = \varphi(x) h(t).$$

Thus

$$\rho(x) \varphi(x) h''(t) = T_0 \varphi''(x) h(t)$$

$$\frac{h''(t)}{h(t)} = \frac{T_0 \varphi''(x)}{\rho(x) \varphi(x)} = -\lambda \quad \leftarrow \text{doesn't depend on } x \text{ or } t.$$

\uparrow only a func of t \uparrow only a func of x

We get 2 ODEs.

$$h''(t) = -\lambda h(t) \quad \text{and} \quad T_0 \varphi''(x) = -\lambda \rho(x) \varphi(x)$$

$\varphi(0) = 0$ and $\varphi(L) = 0$

Compare with

PDE

$$\frac{d}{dx} \left(p(x) \frac{d\phi}{dx} \right) + q(x)\phi + \lambda\sigma(x)\phi = 0,$$

↙ eigenvalue

$$a < x < b,$$

$$L = \frac{d}{dx} p(x) \frac{d}{dx} + q(x)$$

$$L\phi = \frac{d}{dx} \left(p(x) \frac{d\phi}{dx} \right) + q(x)\phi$$

BC

$$\beta_1\phi(a) + \beta_2 \frac{d\phi}{dx}(a) = 0$$

$$\beta_3\phi(b) + \beta_4 \frac{d\phi}{dx}(b) = 0,$$

$$\text{and } p > 0 \text{ and } \sigma > 0$$

$$T_0 \phi''(x) = -\lambda \rho(x) \phi(x)$$

$$\phi(0) = 0 \text{ and } \phi(L) = 0$$

$$a=0 \quad b=L$$

$$\beta_1=1, \beta_2=0$$

$$\beta_3=1, \beta_4=0$$

So this ODE obtained from separation of variables is a Sturm-Liouville problem... Therefore all of the following holds:

1. All the eigenvalues λ are real.
2. There exist an infinite number of eigenvalues:

$$\lambda_1 < \lambda_2 < \dots < \lambda_n < \lambda_{n+1} < \dots$$
 - a. There is a smallest eigenvalue, usually denoted λ_1 .
 - b. There is not a largest eigenvalue and $\lambda_n \rightarrow \infty$ as $n \rightarrow \infty$.
3. Corresponding to each eigenvalue λ_n , there is an eigenfunction, denoted $\phi_n(x)$ (which is unique to within an arbitrary multiplicative constant). $\phi_n(x)$ has exactly $n - 1$ zeros for $a < x < b$.
4. The eigenfunctions $\phi_n(x)$ form a "complete" set, meaning that any piecewise smooth function $f(x)$ can be represented by a generalized Fourier series of the eigenfunctions:

$$f(x) \sim \sum_{n=1}^{\infty} a_n \phi_n(x).$$

Furthermore, this infinite series converges to $[f(x+) + f(x-)]/2$ for $a < x < b$ (if the coefficients a_n are properly chosen).

5. Eigenfunctions belonging to different eigenvalues are orthogonal relative to the weight function $\sigma(x)$. In other words,

$$\int_a^b \phi_n(x) \phi_m(x) \sigma(x) dx = 0 \quad \text{if } \lambda_n \neq \lambda_m.$$

6. Any eigenvalue can be related to its eigenfunction by the Rayleigh quotient:

$$\lambda = \frac{-p\phi \frac{d\phi}{dx} \Big|_a^b + \int_a^b [p(d\phi/dx)^2 - q\phi^2] dx}{\int_a^b \phi^2 \sigma dx},$$

where the boundary conditions may somewhat simplify this expression.

Let's look at the Rayleigh quotient

$$\rho(x) = T_0 > 0$$

$$\sigma(x) = \rho(x) > 0$$

$$q(x) = 0$$

$$\lambda = \frac{-p\phi \frac{d\phi}{dx} \Big|_a^b + \int_a^b [p(d\phi/dx)^2 - q\phi^2] dx}{\int_a^b \phi^2 \sigma dx}$$

boundary conditions may somewhat simplify

$$\phi(0) = 0 \quad \text{and} \quad \phi(L) = 0$$

$$\lambda = \frac{-T_0 \phi(x) \phi'(x) \Big|_0^L + \int_0^L T_0^2 (\phi'(x))^2 dx}{\int_0^L (\phi(x))^2 \rho(x) dx} = \frac{\int_0^L T_0^2 (\phi'(x))^2 dx}{\int_0^L (\phi(x))^2 \rho(x) dx}$$

Note that the integrals are always positive, so $\lambda > 0$.

Since $\lambda > 0$ we can now solve

$$h''(t) = -\lambda h(t)$$

in general as

$$h(t) = c_1 \cos \sqrt{\lambda} t + c_2 \sin \sqrt{\lambda} t \quad \dots \text{ where } \lambda \text{ is an eigenvalue of the Sturm-Liouville problem for } \phi.$$

4. The eigenfunctions $\phi_n(x)$ form a "complete" set, meaning that any piecewise smooth function $f(x)$ can be represented by a generalized Fourier series of the eigenfunctions:

$$f(x) \sim \sum_{n=1}^{\infty} a_n \phi_n(x).$$

Furthermore, this infinite series converges to $[f(x+) + f(x-)]/2$ for $a < x < b$ (if the coefficients a_n are properly chosen).

Let λ_n, ϕ_n be the eigenvalue-eigenfunction pairs corresponding to the Sturm-Liouville problem, and by superposition write

$$u(x,t) = \sum_{n=1}^{\infty} a_n (\cos \sqrt{\lambda_n} t) \phi_n(x) + b_n (\sin \sqrt{\lambda_n} t) \phi_n(x)$$

Now solve for the a_n and b_n in terms of

IC:
$$\begin{aligned} u(x, 0) &= f(x) \\ \frac{\partial u}{\partial t}(x, 0) &= g(x). \end{aligned}$$

$$u(x, 0) = \sum_{n=1}^{\infty} a_n \phi_n(x) = f(x)$$

Eigenfunctions belonging to different eigenvalues are orthogonal relative to the weight function $\sigma(x)$. In other words,

$$\int_a^b \phi_n(x) \phi_m(x) \sigma(x) dx = 0 \quad \text{if } \lambda_n \neq \lambda_m.$$

Any eigenvalue can be related to its eigenfunction by the

Need to add the weight $\sigma(x) = \rho(x)$ to use orthogonality to solve for the coefficients. Therefore,

$$\sum_{n=1}^{\infty} a_n \phi_n(x) \phi_m(x) \rho(x) = f(x) \phi_m(x) \rho(x)$$

Now integrate

$$\int_0^L \sum_{n=1}^{\infty} a_n \phi_n(x) \phi_m(x) \rho(x) dx = \int_0^L f(x) \phi_m(x) \rho(x) dx$$

$$\sum_{n=1}^{\infty} a_n \int_0^L \phi_n(x) \phi_m(x) \rho(x) dx = \int_0^L f(x) \phi_m(x) \rho(x) dx$$

≈ 0 when $n \neq m$

$$\frac{\int_0^L f(x) \phi_n(x) \rho(x) dx}{\int_0^L \phi_n^2(x) \rho(x) dx} = a_n \quad a_m \int_0^L (\phi_m(x))^2 \rho(x) dx = \int_0^L f(x) \phi_m(x) \rho(x) dx$$

Thus

$$(*) \quad a_m = \frac{\int_0^L f(x) \phi_m(x) \rho(x) dx}{\int_0^L (\phi_m(x))^2 \rho(x) dx}$$

Now solve for the other initial condition and the b_n 's.

$$u(x,t) = \sum_{n=1}^{\infty} a_n (\cos \sqrt{\lambda_n} t) \phi_n(x) + b_n (\sin \sqrt{\lambda_n} t) \phi_n(x)$$

$$\frac{\partial u}{\partial t}(x,0) = g(x).$$

$$u_t(x,t) = \sum_{n=1}^{\infty} -a_n \sqrt{\lambda_n} (\sin \sqrt{\lambda_n} t) \phi_n(x) + b_n \sqrt{\lambda_n} (\cos \sqrt{\lambda_n} t) \phi_n(x)$$

$$u_t(x,0) = \sum_{n=1}^{\infty} b_n \sqrt{\lambda_n} \phi_n(x) = g(x)$$

Multiply by the weight and use orthogonality

$$\int_0^L \sum_{n=1}^{\infty} b_n \sqrt{\lambda_n} \phi_n(x) \phi_m(x) \rho(x) dx = \int_0^L g(x) \phi_m(x) \rho(x) dx$$

$$b_m \sqrt{\lambda_m} \int_0^L (\phi_m(x))^2 \rho(x) dx = \int_0^L g(x) \phi_m(x) \rho(x) dx$$

Therefore

$$(**) \quad b_m = \frac{\int_0^L g(x) \phi_m(x) \rho(x) dx}{\sqrt{\lambda_m} \int_0^L (\phi_m(x))^2 \rho(x) dx}$$

$$b_n \sqrt{\lambda_n} = \frac{\int_0^L g(x) \phi_n(x) \rho(x) dx}{\int_0^L \phi_n^2 \rho dx}$$

Solution is

$$u(x,t) = \sum_{n=1}^{\infty} a_n (\cos \sqrt{\lambda_n} t) \phi_n(x) + b_n (\sin \sqrt{\lambda_n} t) \phi_n(x)$$

where the a_n and b_n are determined by (*) and (**) and the λ_n, ϕ_n come from the Sturm-Liouville problem.