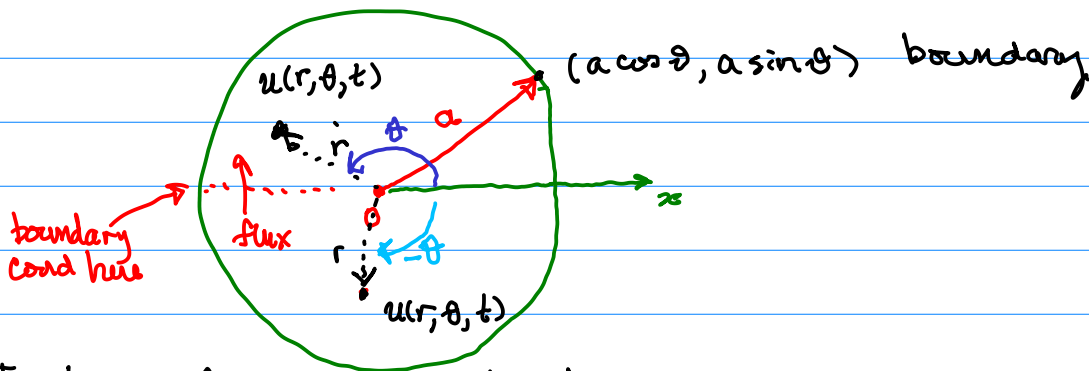


Recall:

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = \frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2}$$

Solving the heat equation $\frac{\partial u}{\partial t} = k \nabla^2 u$ on the disk. u is the temp.



In terms of polar coordinates

PDE $\frac{\partial u}{\partial t} = k \left(\frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} \right)$ $\theta \in [-\pi, \pi]$
 $r \in [0, a]$
 $t \geq 0$

IC $u(r, \theta, 0) = u_0(r, \theta)$

BC Heat bath on boundary $u(a, \theta, t) = f(\theta, t)$
 $u(r, -\pi, t) = u(r, \pi, t)$
 $\frac{\partial u}{\partial \theta}(r, -\pi, t) = \frac{\partial u}{\partial \theta}(r, \pi, t)$

at $r=0$ is where the θ dependency is gone we will assume u is cont. at $r=0$.

Look for equilibrium solutions... No time... boundary conditions don't depend on time... then as $t \rightarrow \infty$ the time dependent solution converges to the equilibrium solution.

Laplace eq. on the disc is equilibrium solution

PDE

$$\frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} = 0$$

$$\theta \in [-\pi, \pi]$$

$$r \in [0, a]$$

B.C.

$$u(a, \theta) = f(\theta)$$

} not homogeneous.

$$u(r, -\pi) = u(r, \pi)$$

$$\frac{\partial u}{\partial \theta}(r, -\pi) = \frac{\partial u}{\partial \theta}(r, \pi)$$

} homogeneous (in θ direction)
(periodic boundary cond.)

Use separation of variables and superposition to solve.

$$u(r, \theta) = \phi(\theta) G(r)$$

Plug it in

orthogonal eigenfunctions.

$$\frac{\partial^2}{\partial r^2} (\phi(\theta) G(r)) + \frac{1}{r} \frac{\partial}{\partial r} (\phi(\theta) G(r)) + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} (\phi(\theta) G(r)) = 0$$

$$\phi(\theta) \frac{\partial^2}{\partial r^2} (G(r)) + \phi(\theta) \frac{1}{r} \frac{\partial}{\partial r} (G(r)) + G(r) \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} (\phi(\theta)) = 0$$

$$G(r) \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} (\phi(\theta)) = -\phi(\theta) \left\{ \frac{\partial^2}{\partial r^2} (G(r)) + \frac{1}{r} \frac{\partial}{\partial r} (G(r)) \right\}$$

$$\frac{1}{\phi(\theta)} \frac{\partial^2}{\partial \theta^2} (\phi(\theta)) = -\frac{r^2}{G(r)} \left\{ \frac{\partial^2}{\partial r^2} (G(r)) + \frac{1}{r} \frac{\partial}{\partial r} (G(r)) \right\} = -\lambda$$

doesn't depend on r doesn't depend on θ \uparrow const.

We obtain two ODEs

$$\frac{\partial^2}{\partial \theta^2} (\varphi(\theta)) = -\lambda \varphi(\theta)$$

$$\varphi(-\pi) = \varphi(\pi)$$

$$\varphi'(-\pi) = \varphi'(\pi)$$

Solve for φ :

Case $\lambda > 0$ $\varphi'' = -\lambda \varphi$

Gen. solution: $\varphi(\theta) = C_1 \cos \sqrt{\lambda} \theta + C_2 \sin \sqrt{\lambda} \theta$

$$\varphi(-\pi) = C_1 \cos \sqrt{\lambda} \pi - C_2 \sin \sqrt{\lambda} \pi$$

$$\varphi(\pi) = C_1 \cos \sqrt{\lambda} \pi + C_2 \sin \sqrt{\lambda} \pi$$

If $\varphi(-\pi) = \varphi(\pi)$ then $2C_2 \sin \sqrt{\lambda} \pi = 0$

$$\varphi'(\theta) = -C_1 \sqrt{\lambda} \sin \sqrt{\lambda} \theta + C_2 \sqrt{\lambda} \cos \sqrt{\lambda} \theta$$

$$\varphi'(-\pi) = C_1 \sqrt{\lambda} \sin \sqrt{\lambda} \pi + C_2 \sqrt{\lambda} \cos \sqrt{\lambda} \pi$$

$$\varphi'(\pi) = -C_1 \sqrt{\lambda} \sin \sqrt{\lambda} \pi + C_2 \sqrt{\lambda} \cos \sqrt{\lambda} \pi$$

If $\varphi'(-\pi) = \varphi'(\pi)$ then $2C_1 \sin \sqrt{\lambda} \pi = 0$

Therefore if $\sin \sqrt{\lambda} \pi \neq 0$ both $C_1 = 0$ and $C_2 = 0$. Therefore we choose $\sqrt{\lambda} = n$ so that $\sin n\pi = 0$.

Case $\lambda = 0$ $\varphi'' = 0$

gen. soln. $\varphi(\theta) = C_1 \theta + C_2$

after boundary cond. we get $C_1 = 0$ so $\varphi(\theta) = C_2$ is another eigen function.

$$r^2 \left\{ \frac{\partial^2}{\partial r^2} (G(r)) + \frac{1}{r} \frac{\partial}{\partial r} (G(r)) \right\} = \lambda G(r)$$

$G(a) =$ by super position

$G(0)$ for continuity need to be bounded
boundary for PDE

$$u(a, \theta) = f(\theta)$$

$$u(r, -\pi) = u(r, \pi)$$

$$\frac{\partial u}{\partial \theta}(r, -\pi) = \frac{\partial u}{\partial \theta}(r, \pi)$$

Next solve the other ODE.

$$r^2 \left\{ \frac{d^2}{dr^2} (G(r)) + \frac{1}{r} \frac{d}{dr} (G(r)) \right\} = \lambda G(r)$$

$$r^2 \frac{d^2}{dr^2} (G(r)) + r \frac{d}{dr} (G(r)) - n^2 G(r) = 0 \quad \text{for } n=0,1,2,\dots$$

This is an Euler equation. The way to solve... plug in $G = r^m$

$$r^2 \frac{d^2}{dr^2} (r^m) + r \frac{d}{dr} r^m - n^2 r^m = 0$$

$$r^2 m(m-1) r^{m-2} + r m r^{m-1} - n^2 r^m = 0$$

Thus $m(m-1) + m - n^2 = 0$

$$m^2 = n^2 \quad \text{so } m = \pm n.$$

General solution $n=1,2,\dots$

$$G(r) = C_1 r^n + C_2 r^{-n}$$

requiring u continuous at $r=0$ means $C_2=0$.

What about $n=0$ (could use reduction of order but easier)

$$r^2 \frac{d^2}{dr^2} (G(r)) + r \frac{d}{dr} (G(r)) = 0$$

$$r \frac{d}{dr} \left(r \frac{d}{dr} G(r) \right) = 0$$

Thus $r \frac{d}{dr} G(r) = C_1 \quad \frac{d}{dr} G(r) = \frac{C_1}{r} \quad G(r) = C_1 \log r + C_2$