

heat equation

$$u_t = k u_{xx}$$

Example of parabolic PDE

time evolution of

a diffusive process.

The heat spreads out

Laplace equation

$$u_{xx} + u_{yy} = 0$$

Example of elliptic, PDE

stationary state of a 2D
heat equation. **No time.**
statics problem

Wave equation

$$u_{tt} = c^2 u_{xx}$$

Example of hyperbolic PDE

time evolution of a
conservation law.

(no dissipation)

CHAPTER 4

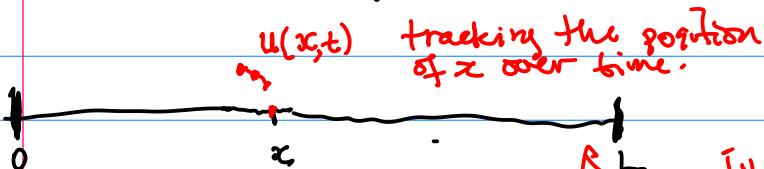
Wave Equation: Vibrating Strings and Membranes

Derivation from Newton's law

$$F = ma$$

mass
force
acceleration.

Material trajectory of that
bit of string marked by x .



equilibrium
position of
a string.

mass	$[M]$	grams.
position	$[L]$	inches cm
velocity	$[L]/[T]$	mph.
acceleration	$[L]/[T]^2$	

examples

position of the point on the string marked by x at time t
is given $u(x, t)$

$$[u] = [L]$$

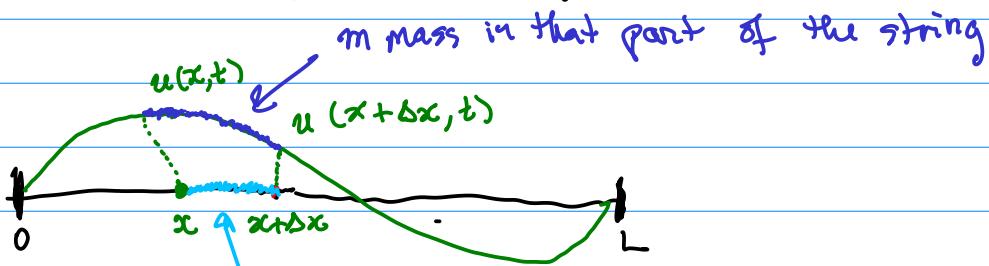
velocity $v = \frac{\partial}{\partial t} u(x, t)$

$$[v] = [L]/[T]$$

acceleration $a = \frac{\partial^2}{\partial t^2} u(x, t)$

$$[a] = [L]/[T]^2$$

$\rho(x)$ the density of the string at the point x

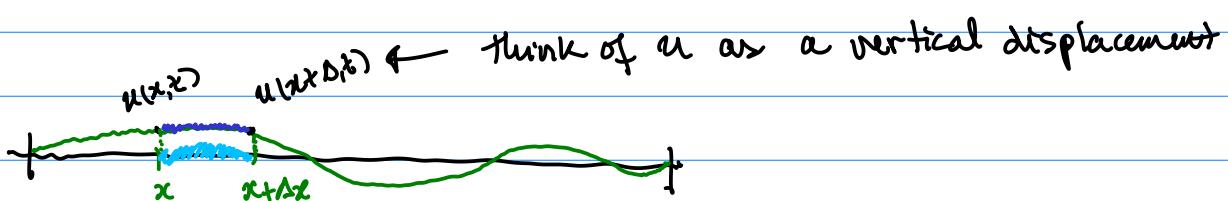


Same as the
mass between
 x and $x + \Delta x$
in the equilibrium position

graph of the position of the
string at a fixed time t

$$m = \int_x^{x+\Delta x} \rho(x) dx$$

Simplifying assumption the displacement from the equilibrium position is small enough the the material trajectories are only in vertical direction



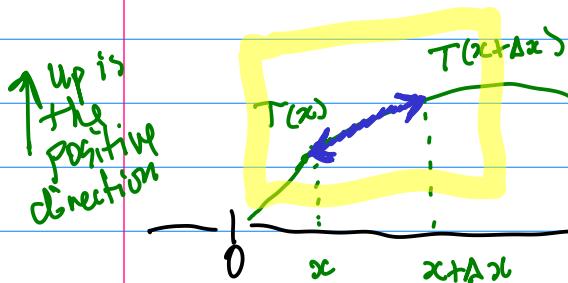
So far

$$F = ma = \left(\int_x^{x+\Delta x} \rho(x) dx \right) u_{tt}(x, t) \approx \rho(x) \Delta x u_{tt}(x, t)$$

we'll take $\Delta x \rightarrow 0$ in just
a moment and in the
limit this is exact.

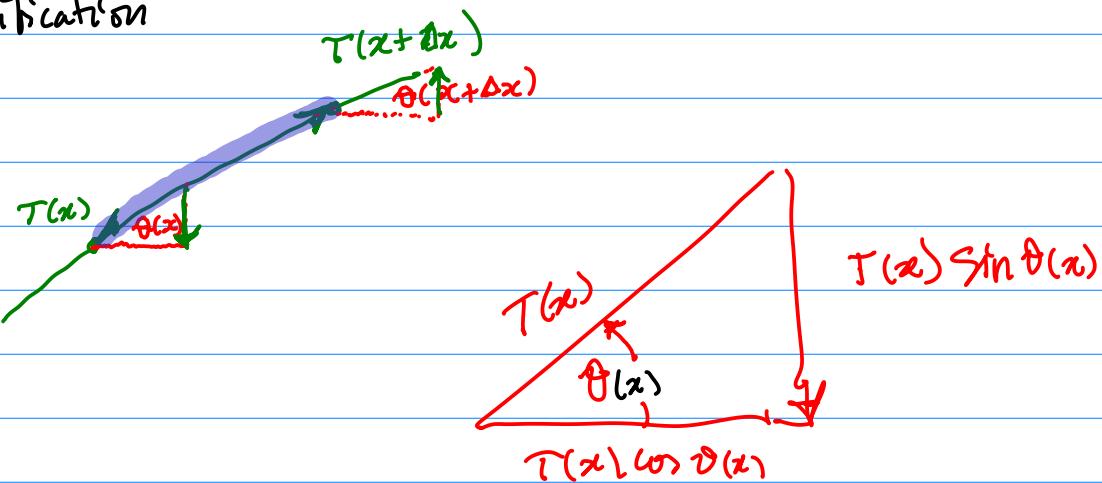
a type of force

Let $T(x)$ be the tension in the string at point x



by assumption there is no horizontal movement so
so don't worry about the force in the horizontal direction.

Magnification



Vertical component of force on the mass between x and $x+\Delta x$

$$f = T(x+\Delta x) \sin \theta(x+\Delta x) \sim T(x) \sin \theta(x)$$

$$F = ma \quad \text{or} \quad m a = F$$

Therefore

$$\rho(x) \Delta x u_{tt}(x, t) = T(x+\Delta x) \sin \theta(x+\Delta x) \sim T(x) \sin \theta(x)$$

Now take the limit as $\Delta x \rightarrow 0$

$$\rho(x) u_{tt}(x, t) = \lim_{\Delta x \rightarrow 0} \frac{T(x+\Delta x) \sin \theta(x+\Delta x)}{\Delta x} \sim T(x) \sin \theta(x)$$

$$\rho(x) u_{tt}(x, t) = \frac{\partial}{\partial x} (T(x) \sin \theta(x))$$

recall definition of derivative

$$f'(x) = \lim_{h \rightarrow 0} \frac{f(x+h) - f(x)}{h}$$

Now take $f(x) = T(x) \sin \theta(x)$
 $h = \Delta x$

$$\frac{1}{\Delta x} \left(\int_x^{x+\Delta x} p(x) dx \right) u_{tt}(x, t) \approx \frac{T(x+\Delta x) \sin \theta(x+\Delta x) - T(x) \sin \theta(x)}{\Delta x}$$

average value
of p on the
interval x to $x+\Delta x$

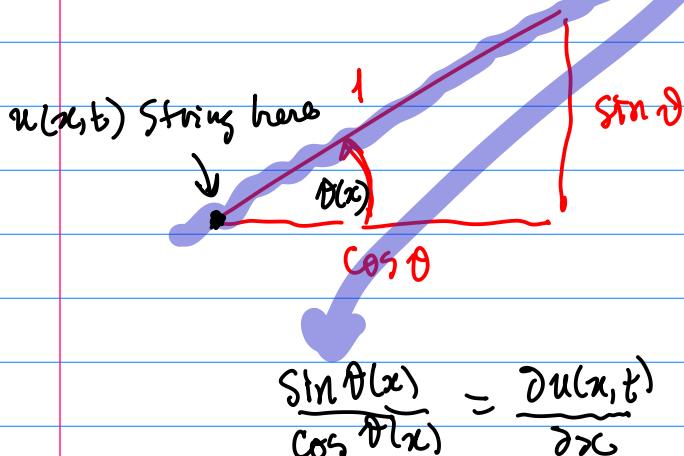
$$\Delta x \rightarrow 0$$

Therefore ...

$$p(x) u_{tt}(x, t) = \frac{\partial}{\partial x} (T(x) \sin \theta(x))$$

write this in terms of u

also might depend on u , but we'll
take it constant for simplicity
instead...



slope of the string is $\frac{\partial u(x, t)}{\partial x}$

also given by $\frac{\sin \theta(x)}{\cos \theta(x)}$

since the displacement is
assumed small we might
as well assume θ is small

$$\sin \theta(x) = \cos \theta(x) \frac{\partial u(x,t)}{\partial x} . \quad \text{then } \cos \theta \approx 1$$

$$\sin \theta(x) \approx \frac{\partial u}{\partial x}(x,t) ,$$

...

$$\rho(x) u_{tt}(x,t) = \frac{\partial}{\partial x} \left(T(x) \frac{\partial u(x,t)}{\partial x} \right)$$

For the first example we assume $\rho(x) = \rho_0$ a constant
and $T(x) = T_0$ is constant.